# Parity-conserved self-consistent CHFB solution 

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We developed a new program for solving the constrained Hartree-Fock-Bogoliubov (CHFB) equation without parity mixing. In this scheme (CHFB5), we require five constraints, one each on the total angularmomentum $I$, proton number $Z_{+}$in the + parity shell $\left(p^{+}\right)$, proton number $Z_{-}$in the - parity shell $\left(p^{-}\right)$, neutron number $N_{+}$in the + parity shell $\left(n^{+}\right)$, and neutron number $N_{-}$in the - parity shell $\left(n^{-}\right)$. As an example, we choose ${ }^{134} \mathrm{Nd}$ with the same parameter set as that adopted in Ref. 1). Here, we solved the full CHFB equation ${ }^{2)}$ including all exchange terms (Fock terms), while Ref. 1). adopts only the Hartree terms. The values of $\left(Z_{+}, Z_{-}, N_{+}, N_{-}\right)$are selected in reference to the usual CHFB solutions with three constraints (CHFB3). The usual CHFB3 solutions show $\left(Z_{+}, Z_{-}, N_{+}, N_{-}\right)=(14.59,17.41,13.87,10.13)$ at $I=$ 0 , while $(14.04,17.96,14.0,10.0)$ at $I=26$. Here, $(Z, N)=(32,24)$ are numbers outside the closed core $(28,50)$. Thus, we select $(14,18,14,10)$ for the CHFB5 equation. The intrinsic difference between CHFB3 and CHFB5 solutions is in the quasi-particle (QP) energies. In Fig. 1, we compare the behavior of the lowest QP energies of $\Lambda$ with its time-reversed energy $\tilde{\Lambda}$ vs. $I$. The equations for $\Lambda$ and $\tilde{\Lambda}$ have been provided in Ref. 2). The degeneracy is lifted by the Coriolis antipairing effect with increasing $I$. Figure 1(A) shows the neutron shell, and (B) the proton shell. In both panels, $\pm$ specifies the $\pm$ shell; the filled symbols express $\tilde{\Lambda}$ and the open symbols $\Lambda$. Those in the abbreviation "with" denote CHFB5 solutions, while the others denote CHFB3 solutions. At low $I$, QP energies by CHFB3 and CHFB5 solutions coincide in the neutron shells (A); however, there is a considerable difference among the $p^{+}$shell (B). The negative value of $\Lambda$ in the $n^{+}$shell is observed at $I=10$ in both CHFB3 and CHFB5 solutions; this indicates the first backbending is caused by the $\mathrm{i}_{13 / 2}$ level in the $n^{+}$shell. There occur negative values of $\Lambda$ in the $n^{+}$and $n^{-}$shells around $I=20$ to 26 , and they correspond to decreasing $\Delta_{n}$, i.e., 0.00021 (CHFB3) and 0.00035 (CHFB5).

## References

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Fig. 1. (A) Lowest QP energies $\Lambda$ and $\tilde{\Lambda}$ in the neutron shell as functions of angular momentum $I$. The red opencircles represent $\Lambda$, and the red filled-circles represent $\tilde{\Lambda}$ in the $n^{+}$shell by the CHFB3 solutions; the orange open-triangles denote $\Lambda$, and the orange filled-triangles represent $\tilde{\Lambda}$ in the $n^{+}$shell by the CHFB5 solutions. The open squares represent $\Lambda$, and the filled squares represent $\tilde{\Lambda}$ in the $n^{-}$shell by the CHFB3 solutions, while the blue open-triangles-down represent $\Lambda$, and the blue filled-triangles-down represent $\tilde{\Lambda}$ in the $n^{-}$shell by the CHFB5 solutions. (B) The lowest QP energies of $\Lambda$ and $\tilde{\Lambda}$ in the proton shell as functions of $I$. The red open-circles represent $\Lambda$, and the red filled-circles represent $\tilde{\Lambda}$ in the $p^{+}$shell by the CHFB3 solutions, while the orange open-triangles represent $\Lambda$, and the orange filled-triangles represent $\tilde{\Lambda}$ in the $p^{+}$shell by the CHFB5 solutions. The open squares represent $\Lambda$, and the filled squares represent $\tilde{\Lambda}$ in the $p^{-}$shell by the CHFB3 solutions, while the blue open-triangles-down represent $\Lambda$ and the blue filled-triangles-down represent $\tilde{\Lambda}$ in the $p^{-}$shell by the CHFB5 solutions.


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