# Isomeric-to-ground state ratio of ${ }^{128} \mathrm{Sn}$ measured by Rare RI Ring 

H. F. Li, ${ }^{* 1, * 2, * 3}$ S. Naimi, ${ }^{* 1}$ D. Nagae, ${ }^{* 1}$ Y. Abe, ${ }^{* 1}$ F. Suzaki, ${ }^{* 1}$ Y. Yamaguchi, ${ }^{* 1}$ M. Wakasugi, ${ }^{* 1}$ S. Omika, ${ }^{* 4}$ K. Inomata, ${ }^{* 4}$ H. Arakawa, ${ }^{* 4}$ S. Hosoi, ${ }^{* 4}$ K. Nishimuro, ${ }^{* 4}$ Y. Inada, ${ }^{* 4}$ T. Kobayashi, ${ }^{* 4}$ D. Kajiki, ${ }^{* 4}$<br>D. Hamakawa, ${ }^{* 4}$ W. B. Dou, ${ }^{* 4}$ T. Yamaguchi, ${ }^{* 4}$ M. Mukai, ${ }^{* 5}$ T. Moriguchi, ${ }^{* 5}$ R. Kagesawa, ${ }^{* 5}$ D. Kamioka, ${ }^{* 5}$ A. Ozawa, ${ }^{* 5}$ S. Suzuki, ${ }^{* 2}$ Z. Ge, ${ }^{* 2}$ C. Y. Fu, ${ }^{* 2}$ Q. Wang, ${ }^{* 2}$ M. Wang, ${ }^{* 2}$ S. Ota, ${ }^{* 6}$ S. Michimasa, ${ }^{* 6}$ N. Kitamura, ${ }^{* 6}$ S. Masuoka, ${ }^{* 6}$ D. S. Ahn, ${ }^{* 1}$ H. Suzuki, ${ }^{* 1}$ N. Fukuda, ${ }^{* 1}$ H. Takeda, ${ }^{* 1}$ Y. Shimizu, ${ }^{* 1}$ Y. A. Litvinov, ${ }^{* 7}$ G. Lorusso, ${ }^{* 8}$ and T. Uesaka ${ }^{* 1}$

The excitation energy of the isomeric state and the isomeric-to-ground state ratios are very important to understand the nuclear structure and reactions. Direct mass measurement can be used for measuring the excitation energy of the long-lived isomeric state and determining the isomeric-to-ground state ratio simultaneously. Rare RI Ring (R3) is an isochronous mass spectrometer in RIBF. The principle of the mass measurement at R3 is described by the following equation:

$$
\begin{equation*}
\frac{m_{1}}{q_{1}}=\frac{m_{0}}{q_{0}} \frac{T_{1}^{c o r r}}{T_{0}}=\frac{m_{0}}{q_{0}} \frac{1}{T_{0}} T_{1} \sqrt{1+\frac{1-\left(\frac{T_{0}}{T_{1}}\right)^{2}}{\left(\frac{m_{0} / q_{0}}{B \rho_{1}} c\right)^{2}}} \tag{1}
\end{equation*}
$$

wherec $\mathrm{T}_{1}$ and $\mathrm{T}_{0}$ are the time-of-flight (TOF) of the nucleus of interest and reference nucleus, respectively, and $B \rho_{1}$ is the magnet rigidity of the nucleus of interest. ${ }^{1)}$ The unknown mass $m_{1}$ is determined relative to the mass of the isochronous reference nucleus $m_{0}$. B $\rho$ tagging is performed at the momentum-dispersive focal plane F5 of BigRIPS by measuring the horizontal position with two parallel-plate avalanche counters (PPACs). The TOF of the nuclei in R3 was measured using the E-MCP detector $^{2)}$ at S0 of SHARAQ and a plastic scintillator placed


Fig. 1. (a), (b) Correlations between $\mathrm{T}_{1}$ and $\mathrm{T}_{1}^{\text {corr }}$, respectively, and the F5 position for ${ }^{128} \mathrm{Sn}$.

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Fig. 2. $\mathrm{T}_{1}^{\text {corr }}$ spectrum of ${ }^{128} \mathrm{Sn}$; two Gaussian functions were used to fit the histogram.
behind R3.
In the autumn of 2018, the first mass measurement campaign was conducted at R3. To measure the mass of ${ }^{125} \mathrm{Ag}$ and ${ }^{124} \mathrm{Pd}, 3$ neighbor isotones, ${ }^{126} \mathrm{Cd},{ }^{127} \mathrm{In}$, and
${ }^{128} \mathrm{Sn}$, were injected in R3 as reference nuclei to determine the mean $\mathrm{B} \rho$ value of $\mathrm{R} 3 .{ }^{3)}$ The first isomeric state of ${ }^{128} \mathrm{Sn}$, the excitation energy and half-life of which are about 2091.5 keV and $6.5 \mathrm{~s},{ }^{5)}$ respectively, was produced and observed during this experiment. The TOF in R3 for each particle was normalized to the same turn numbers ${ }^{4)}$ to determine $T_{1}$. The correlation between the $T_{1}$ of ${ }^{128} \mathrm{Sn}$ and the F5 position is shown in Fig. 1(a). After event-wise correction with $\mathrm{B} \rho, \mathrm{T}_{1}^{\text {corr }}$, s dependence on the F5 position is shown in Fig. 1(b). The isomeric state and ground state of ${ }^{128} \mathrm{Sn}$ can be well resolved in the spectrum of $\mathrm{T}_{1}^{\text {corr }}$, as shown in Fig. 2. The left peak is the ground state of ${ }^{128} \mathrm{Sn}$, and the right one is the isomeric state. A function composed of the sum of two Gaussian functions was used to fit this spectrum.

The mass resolving power of R3 can be given by

$$
\begin{equation*}
R=m / \Delta m=T_{1}^{\text {corr }} / \Delta T_{1}^{c o r r}, \tag{2}
\end{equation*}
$$

which is derived from Eq. (1). The full width at half maximum of the achieved mass resolving power is about 125,000 from the ground-state peak. The isomeric-toground state ratio for ${ }^{128} \mathrm{Sn}$ is $1.8(4)$, which is determined by the integral values' ratio of the fitting functions.

## References

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[^0]:    *1 RIKEN Nishina Center
    *2 Institute of Modern Physics, Chinese Academy of Sciences
    *3 University of Chinese Academy of Sciences
    *4 Department of Physics, Saitama University
    *5 Institute of Physics, University of Tsukuba
    *6 Center for Nuclear Study, University of Tokyo
    *7 GSI Helmholtzzentrum für Schwerionenforschung
    *8 NPL, University of Surrey

